

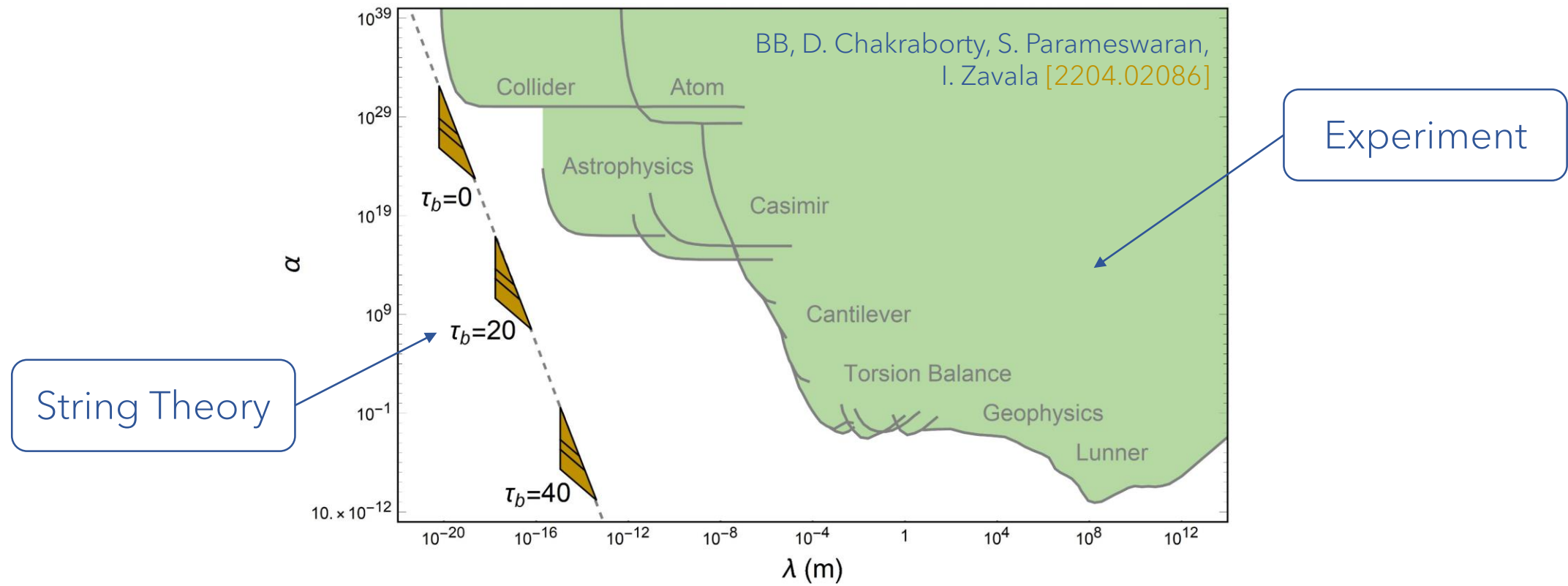
Gravity at the Tip of the Throat

Bruno Valeixo Bento

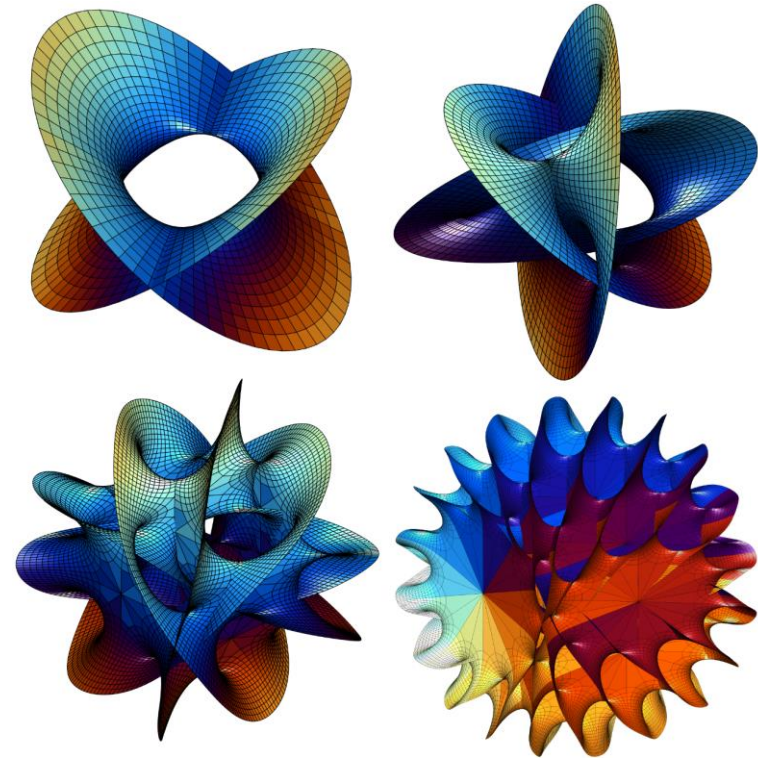
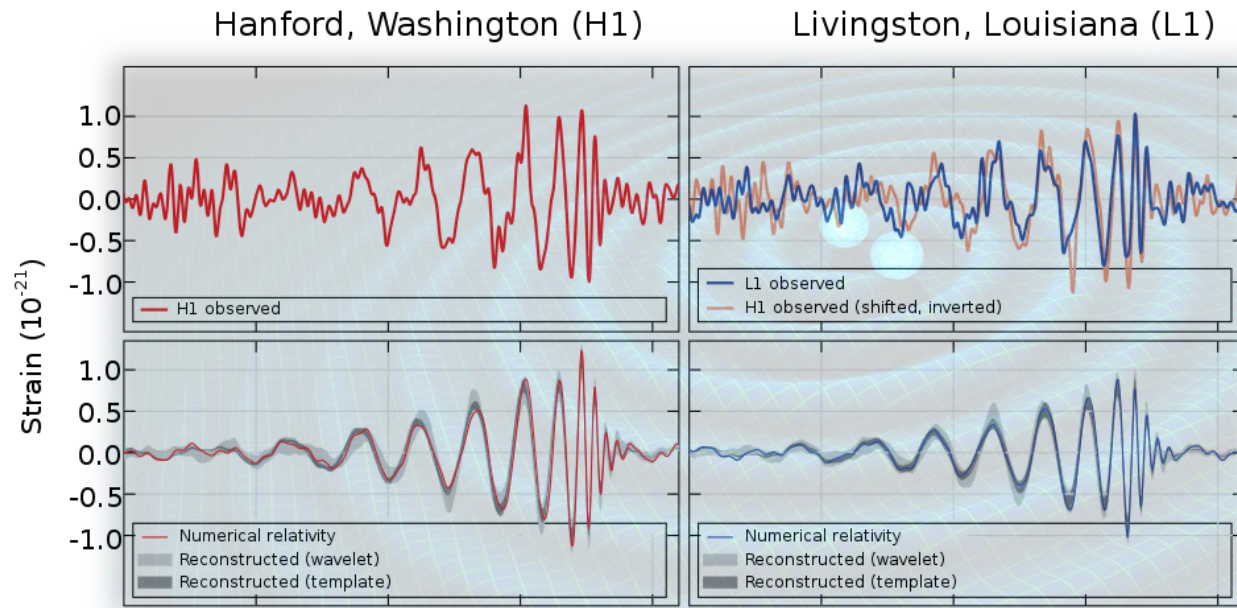
Based on [\[2204.02086\]](#) in collaboration with:

Dibya Chakraborty, Susha Parameswaran, Ivonne Zavala

Spoiler

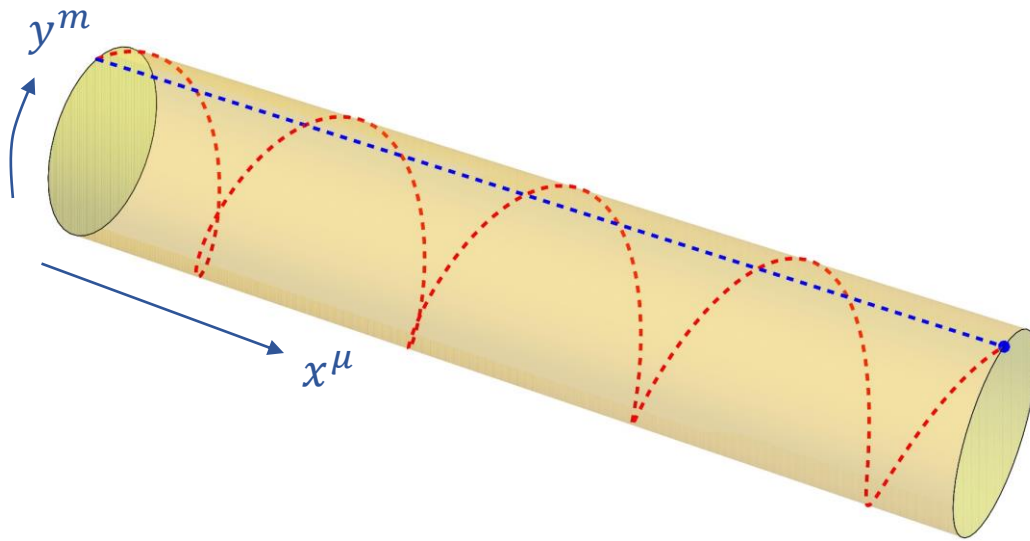


The Question



Can GWs tell us about (warped) extra dimensions?

Kaluza-Klein modes

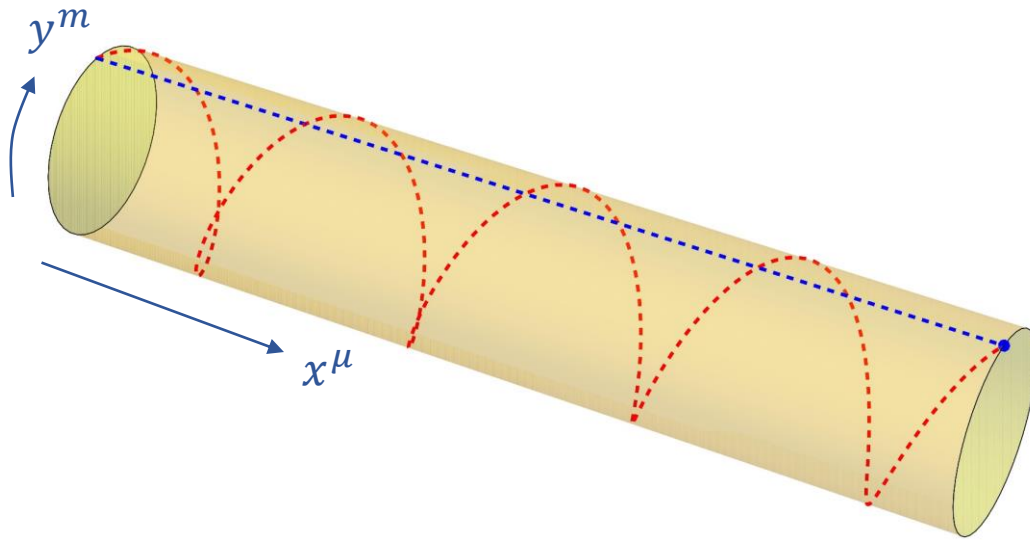


$$p^\mu p_\mu = -m^2 - \underbrace{p^n p_n}_{m_k^2}$$

Diagram illustrating the mass spectrum of Kaluza-Klein modes. The vertical axis represents mass squared, with levels labeled m_4 , m_3 , m_2 , and M_{KK} . The equation $p^\mu p_\mu = -m^2 - p^n p_n$ is shown, with the term $p^n p_n$ identified as m_k^2 .

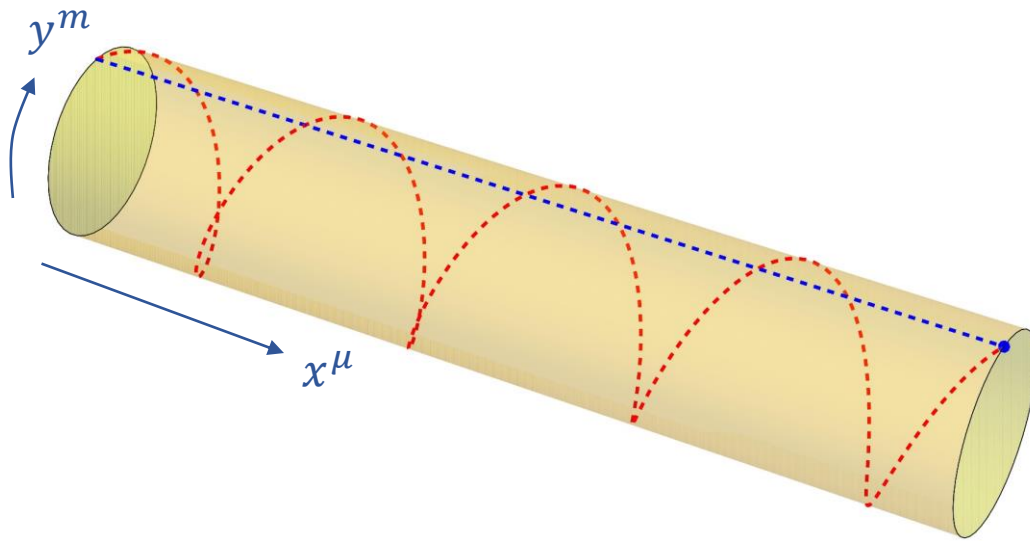
The tower is the signature of **extra dimensions**.

Kaluza-Klein modes



The tower is the signature of **extra dimensions**.

Kaluza-Klein modes



The tower is the signature of **extra dimensions**.

Warped compactification

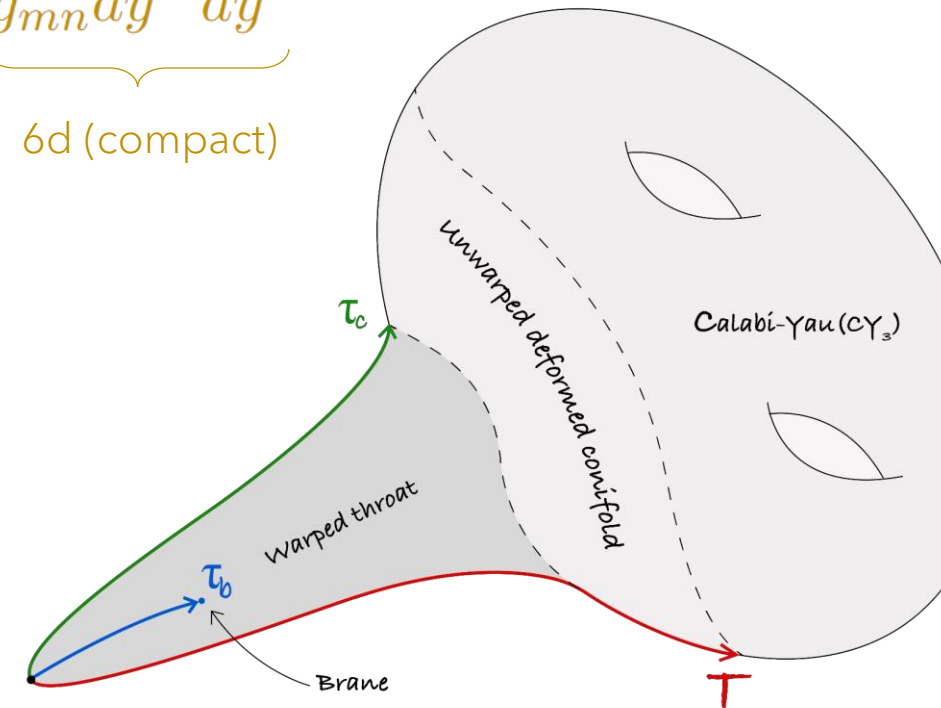
String Theory: 10d \rightarrow 4d

$$ds^2 = \underbrace{H(y)^{-1/2} g_{\mu\nu} dx^\mu dx^\nu}_{4d} + \underbrace{H(y)^{1/2} \tilde{g}_{mn} dy^m dy^n}_{6d \text{ (compact)}}$$

Warp factor

The warping makes scales position (y^m) dependent

What does it do to Gravitational Waves?



KK Gravitational Waves

$$S_{\text{IIB}} = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{-G} (R + \mathcal{L}_{\text{IIB}}) + S_{\text{brane}}$$



$$\square_{10} h_{MN} + 2\bar{R}_{MPNQ} h^{PQ} = T_{MN}^{(1)}$$



$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu}$$

10d gravity (e.g. Type IIB)



10d GWs



4d GWs (tower)

KK Gravitational Waves

$$S_{\text{IIB}} = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{-g} \left(R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \dots \right) \quad (\text{e.g. Type IIB})$$

$$h_{\mu\nu}(x, y) = \sum_k h_{\mu\nu}^k(x) \phi_k(y)$$

wavefunctions

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu}$$

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu}$$

4d GWs (tower)

$$h_{MN} \rightarrow h_{\mu\nu}, h_{\mu n}, h_{mn}$$

KK Gravitational Waves

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu} \quad \leftrightarrow \quad \text{Tower of GWs } (f_{GW} \sim m_k)$$

OBSERVATION

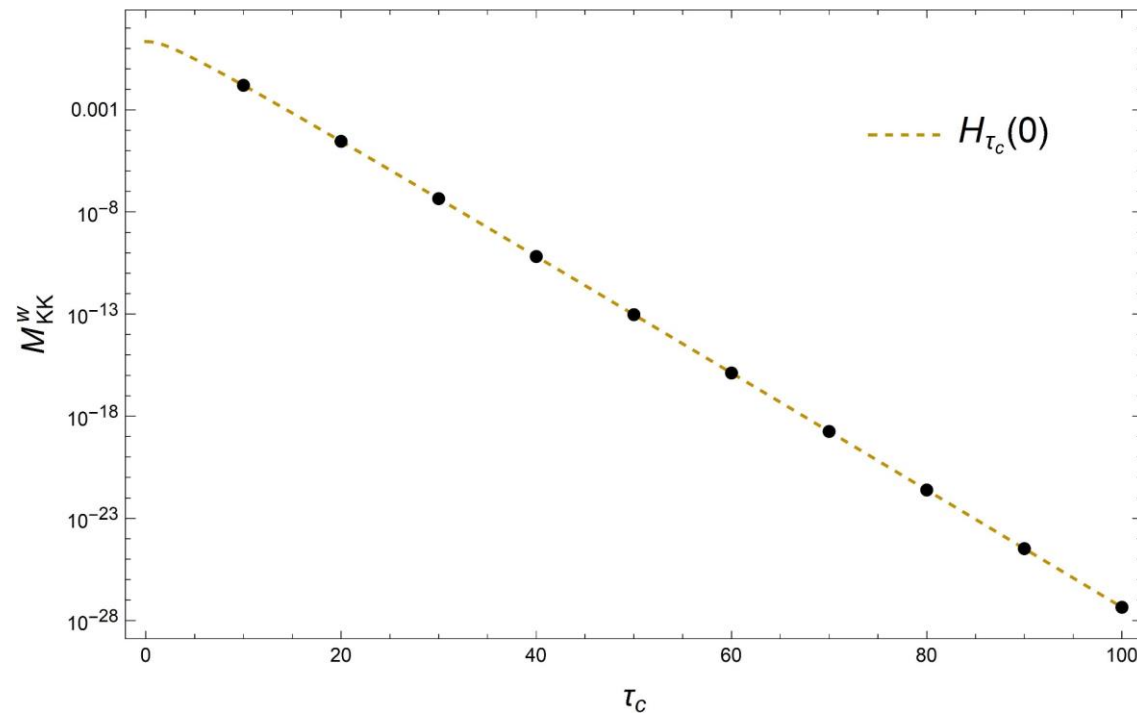
f_{GW} cannot* be too high ($10^4 \text{ Hz} \sim 10^{-23} \text{ eV}$)

QUESTION

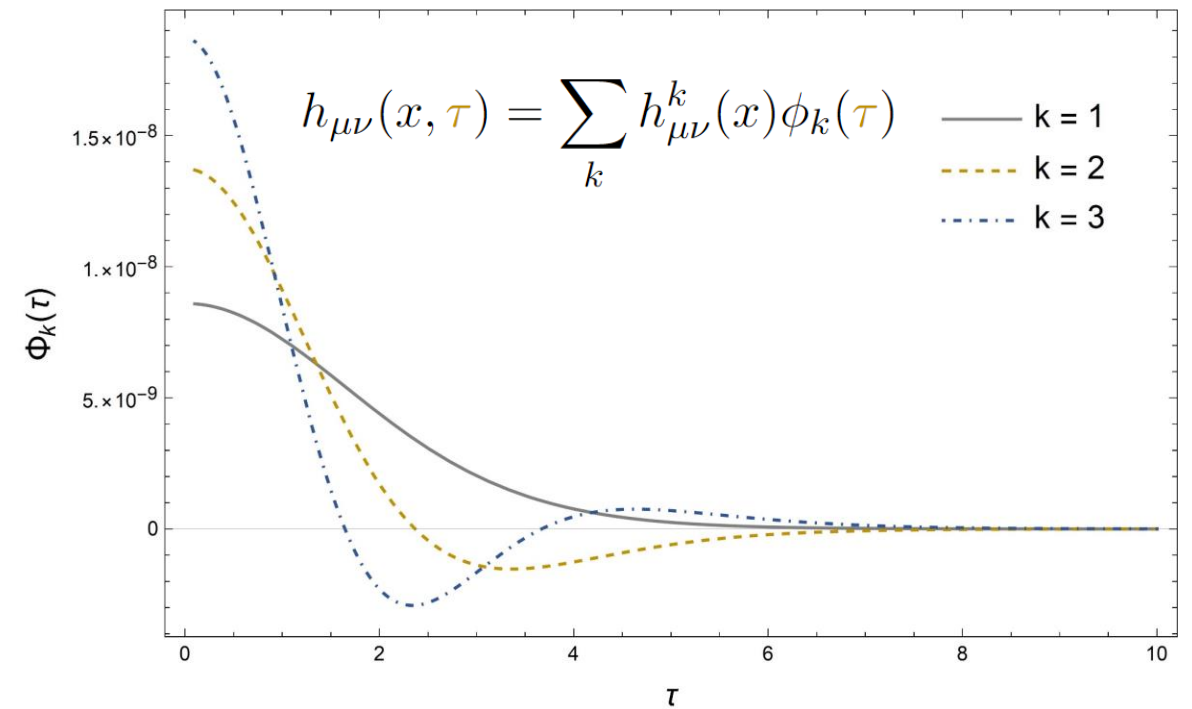
Wouldn't we have seen $h_{\mu\nu}^k$ already?

KK Gravitons

*Fully warped limit ($\tau_c = T$)



Masses



Wavefunctions

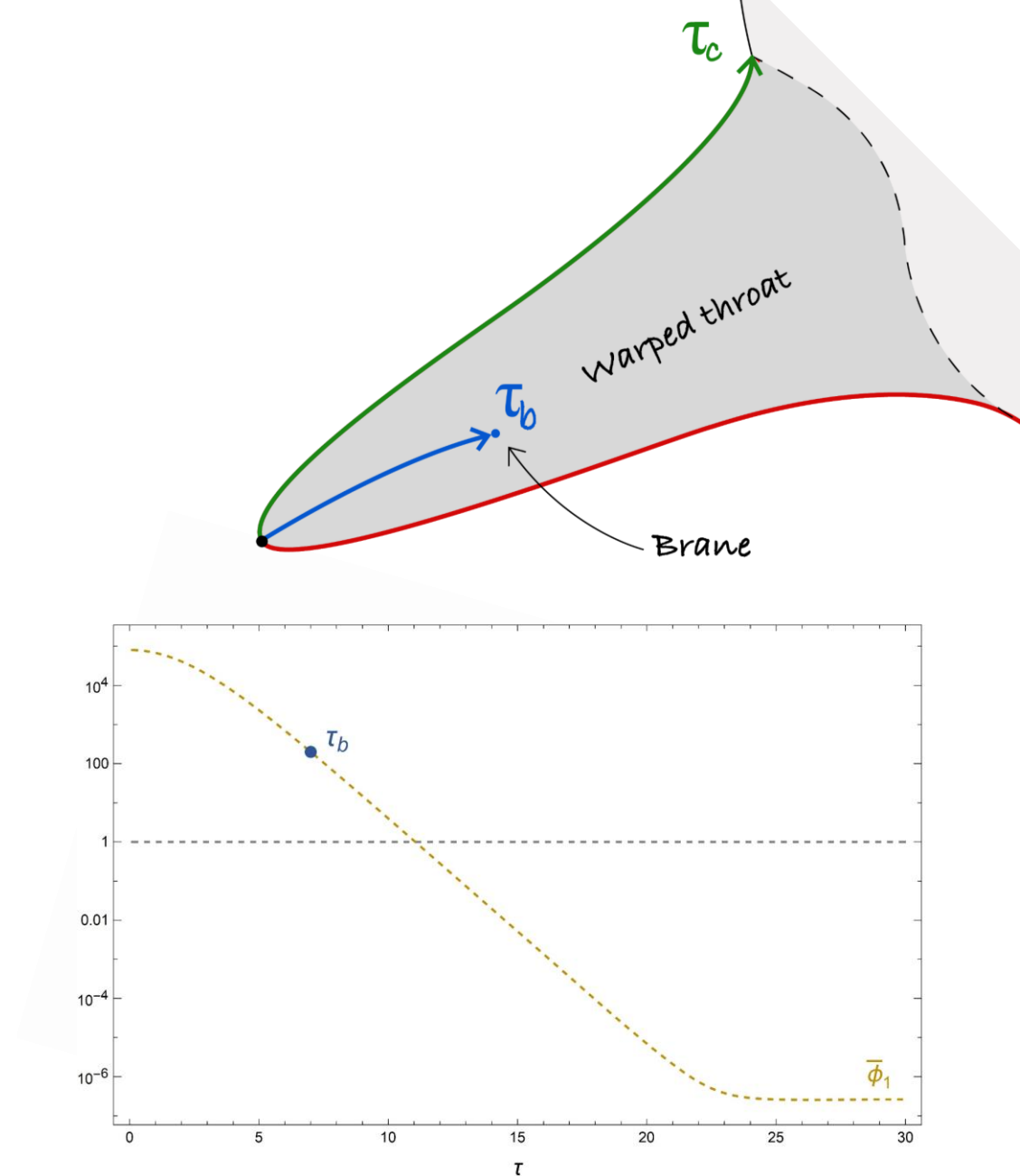
Braneworld

For a **brane** somewhere in the throat

$$\mathcal{L}_{\text{brane}} \sim \sum_k \underbrace{\frac{\bar{\phi}_k(\tau_b)}{M_P}}_{\text{Couplings}} h_{\mu\nu}^k T^{\mu\nu}$$

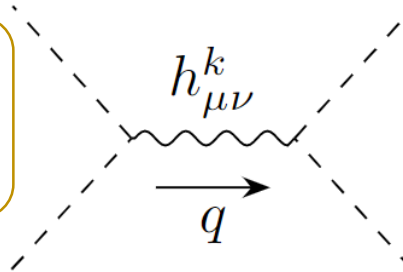
Gravitational interactions on the brane include the **whole tower**

$$h_{\mu\nu}(x, y) = \sum_k h_{\mu\nu}^k(x) \phi_k(y)$$



Do we see the extra dimensions?

The wavefunctions weigh the contribution of each mode

$$V(q) = \lim_{q^0 \rightarrow 0} \sum_k \left| \frac{\bar{\phi}_k(\tau_b)}{M_P} \right|^2$$


The diagram shows a wavy line representing a graviton exchange between two vertices, enclosed in dashed lines. The wavy line is labeled $h_{\mu\nu}^k$ and has a momentum vector q pointing to the right.

To compare with experiments we express it as

$$V(r) = G \frac{m_1 m_2}{r} \left(1 + \alpha e^{-r/\lambda} \right)$$

Strength Range

Do we see the extra dimensions?

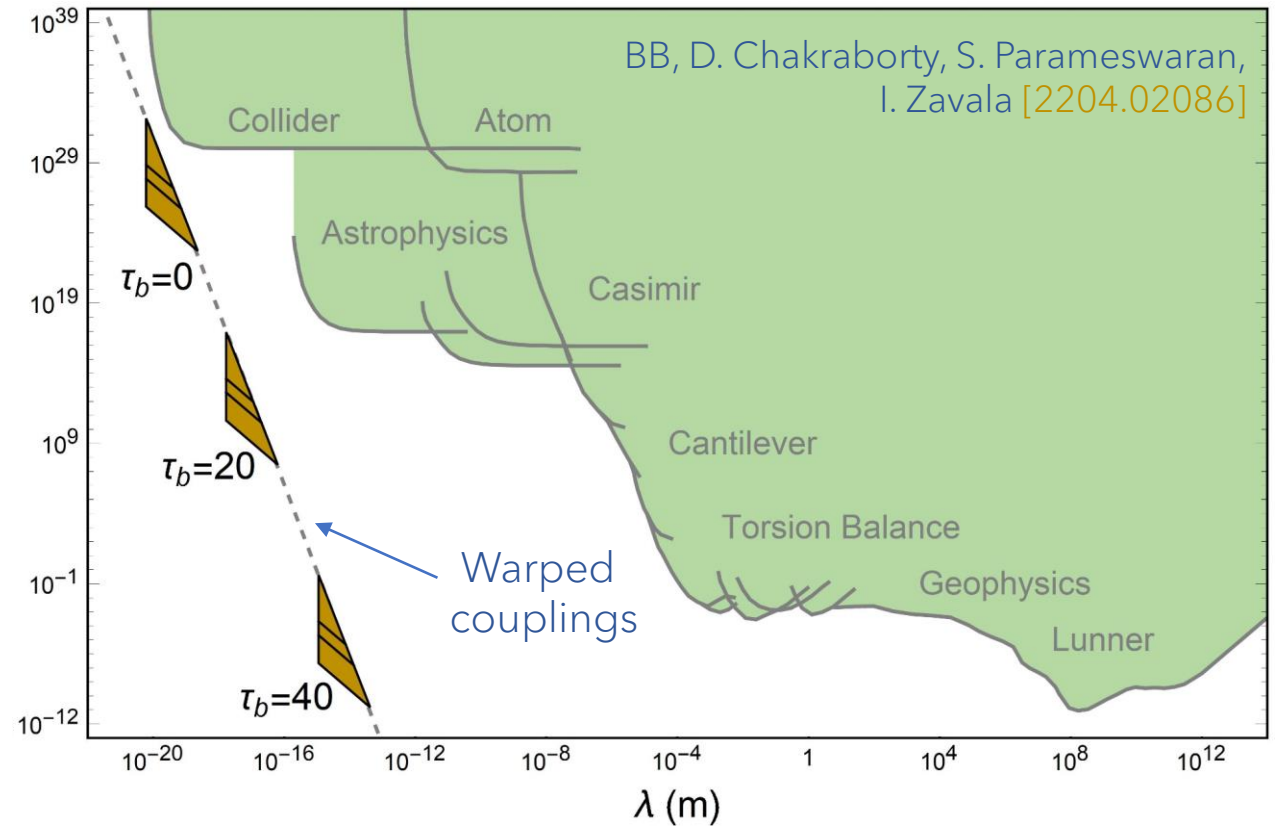
Predictions vs experimental constraints

$$|\bar{\phi}_k|^2 \sim \alpha \approx \frac{(2\pi)^2}{(g_s M)^3} \frac{2|\phi_1(\tau_b)|^2}{I(\tau_b)^{1/2}} \frac{g_s^2}{\mathcal{H}^2}$$

$$m_k \sim \lambda^{-1} \approx \frac{\mathcal{H}}{2^{1/6} \sqrt{g_s M}} I(\tau_b)^{1/4} \quad \alpha$$

Triangle regions:

- $g_s M > 1$ Supergravity (α')
- $g_s < 1$ String loop expansion
- $M < M_{max}$ D3 Tadpole



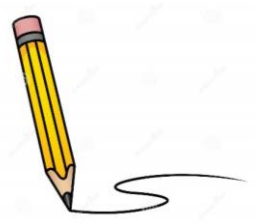
Conclusions

- (compact) Extra dimensions \rightarrow Tower of 4d gravitons \rightarrow Tower of 4d GWs
- Phenomenological parameters (α, λ) \leftrightarrow String theory parameters (g_s, M, \dots)
No conflict
- Wavefunction profiles are important! What's the effect on GWs?

Thank You

BACKUP SLIDES

KK Gravitational Waves



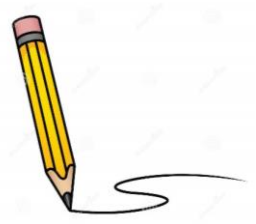
Higher dimensional GWs

$$\square_{10} h_{MN} - 2\bar{R}^S_{MNP} g^{PQ} h_{QS} = T_{MN}^{(1)}$$

give rise 4d GWs

$$\square_4 h_{\mu\nu} + \underbrace{\left(\frac{\Delta_6}{c^{1/2} H} - \frac{2}{3} \Lambda_4 \right)}_{\text{6d (compact)}} h_{\mu\nu} = T_{\mu\nu}^{(1)}$$

$$h_{\mu\nu}(x, y) = \sum_k h_{\mu\nu}^k(x) \phi^k(y)$$



KK Gravitational Waves

Higher dimensional GWs

$$\square_{10} h_{MN} - 2$$

give rise 4d GWs

$$\square_4 h_{\mu\nu} + \left(-\frac{2}{3} \Lambda_4 \right) h_{\mu\nu} = -m_k^2 \phi^k(y)$$

Define the functions as

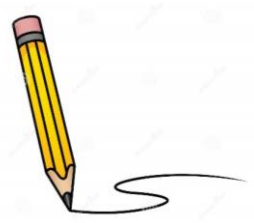
$$\left(\frac{\Delta_6}{c^{1/2} H} - \frac{2}{3} \Lambda_4 \right) \phi^k(y) = -m_k^2 \phi^k(y)$$

\Rightarrow Infinitely many orthogonal modes.

6d (compact)

$$h_{\mu\nu}(x, y) = \sum_k h_{\mu\nu}^k(x) \phi^k(y)$$

KK Gravitational Waves



The decomposition is defined by

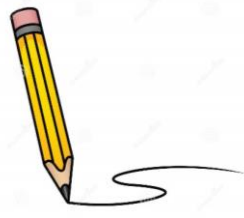
$$\left(\frac{\Delta_6}{c^{1/2}H} - \frac{2}{3}\Lambda_4 \right) \phi^k(y) = -m_k^2 \phi^k(y)$$

i.e. the functions $\phi^k(y)$ are eigenfunctions with eigenvalue $-m_k^2$ and orthogonal

$$\int d^6y \sqrt{g_6} H(y) \phi_k(y) \phi_{k'}(y) = \mathcal{N}_{(k)}^2 \delta_{kk'}$$

Infinitely many solutions (discrete) \Rightarrow Tower of states

KK Gravitational Waves



$$\square_4 h_{\mu\nu} + \left(\frac{\Delta_6}{c^{1/2} H} - \frac{2}{3} \Lambda_4 \right) h_{\mu\nu} = 2\tilde{T}_{\mu\nu}^{(1)}$$

$$h_{\mu\nu}(x, \mathbf{y}) = \sum_{k'} h_{\mu\nu}^{k'}(x) \phi_{k'}(\mathbf{y})$$

$$\sum_{k'} \left(\square_4 h_{\mu\nu}^{k'} - m_{k'}^2 h_{\mu\nu}^{k'} \right) \phi_{k'}(\mathbf{y}) = 2\tilde{T}_{\mu\nu}^{(1)}(x, \mathbf{y})$$

$$\int d^6 \mathbf{y} \sqrt{g_6} H(\mathbf{y}) \phi_k(\mathbf{y})$$

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu}$$

$$T_{\mu\nu} \equiv 2 \int d^6 \mathbf{y} \sqrt{g_6} H(\mathbf{y}) \phi_k(\mathbf{y}) \tilde{T}_{\mu\nu}^{(1)}(x, \mathbf{y})$$

KK Gravitational Waves

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu} \qquad T_{\mu\nu} \equiv 2 \int d^6 y \sqrt{g_6} H(y) \phi_k(y) \tilde{T}_{\mu\nu}^{(1)}(x, y)$$

The warped extra dimensions contribute with

- Tower of GWs ($k > 0$) (high frequencies $f_{GW} \sim m_k$)
- More sources
 - Bulk (e.g. h_{mn})
 - Localised (e.g. brane fields)
- Warped energy-momentum

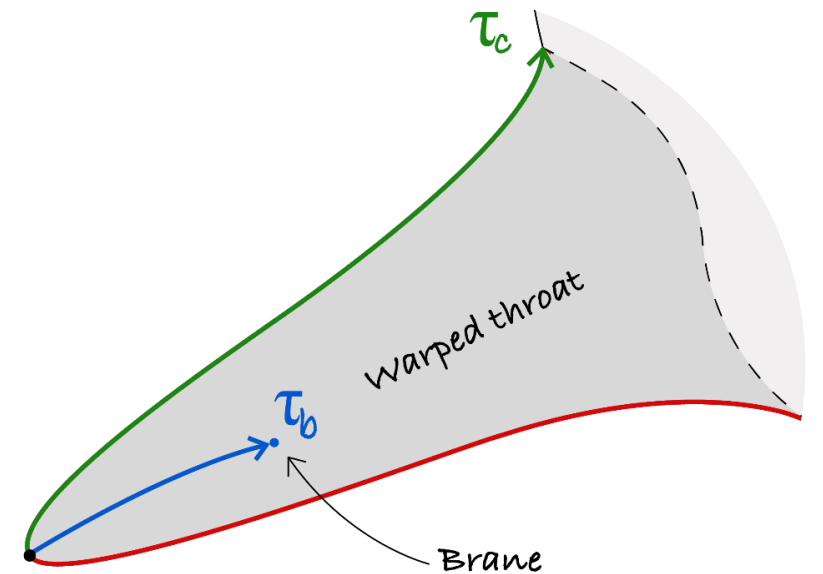
Warped deformed conifold

Klebanov-Strassler (KS) solution

$$ds^2 = H(y)^{-1/2} g_{\mu\nu} dx^\mu dx^\nu + H(y)^{1/2} c^{1/2} g_{mn} dy^m dy^n$$

$$g_{mn} = \begin{cases} g_{mn}^{(\text{KS})} & H(y) \gg 1 \\ g_{mn}^{(\text{CY})} & H(y) \sim 1 \end{cases}$$

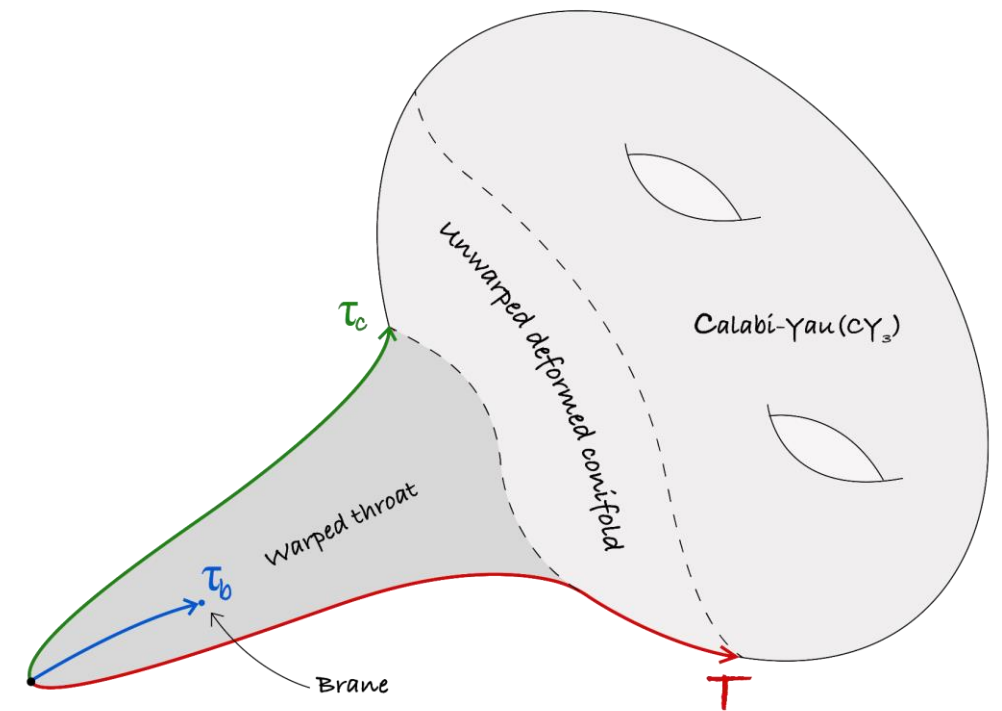
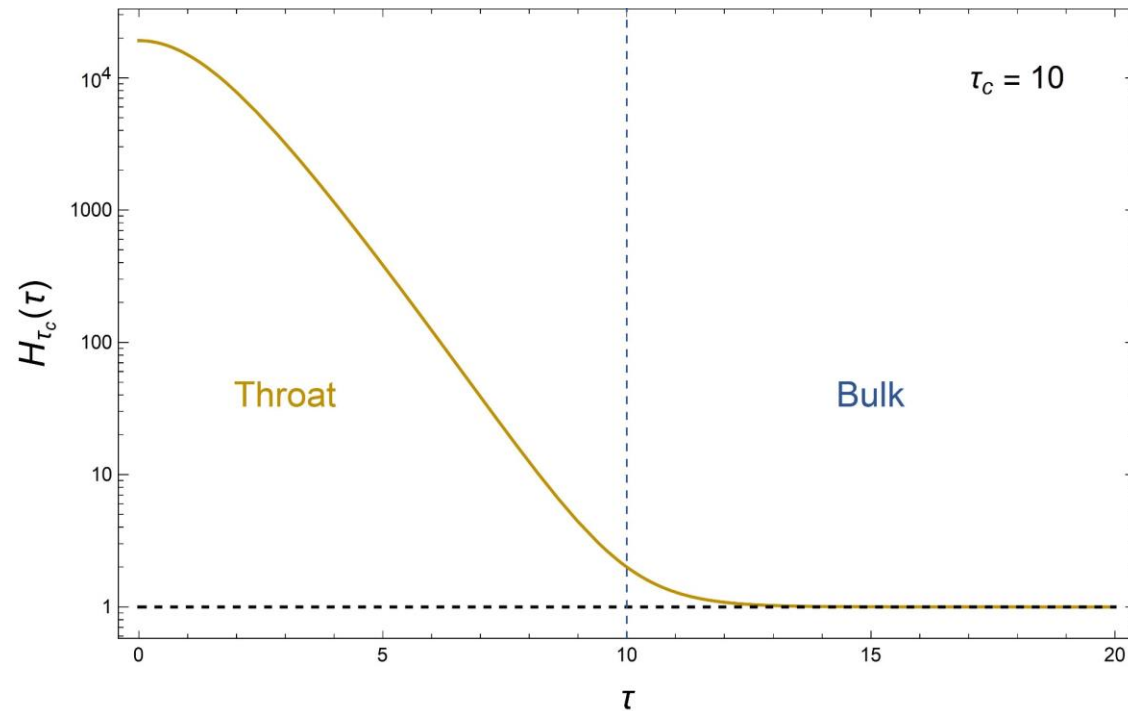
$$g_{mn}^{(KS)} = \frac{\epsilon^{4/3}}{2} \mathcal{K}(\tau) \begin{pmatrix} \frac{1}{3\mathcal{K}^3(\tau)} \mathbb{1}_2 & 0 & 0 \\ 0 & \sinh^2(\tau/2) \mathbb{1}_2 & 0 \\ 0 & 0 & \cosh^2(\tau/2) \mathbb{1}_2 \end{pmatrix}$$



Warped compactification

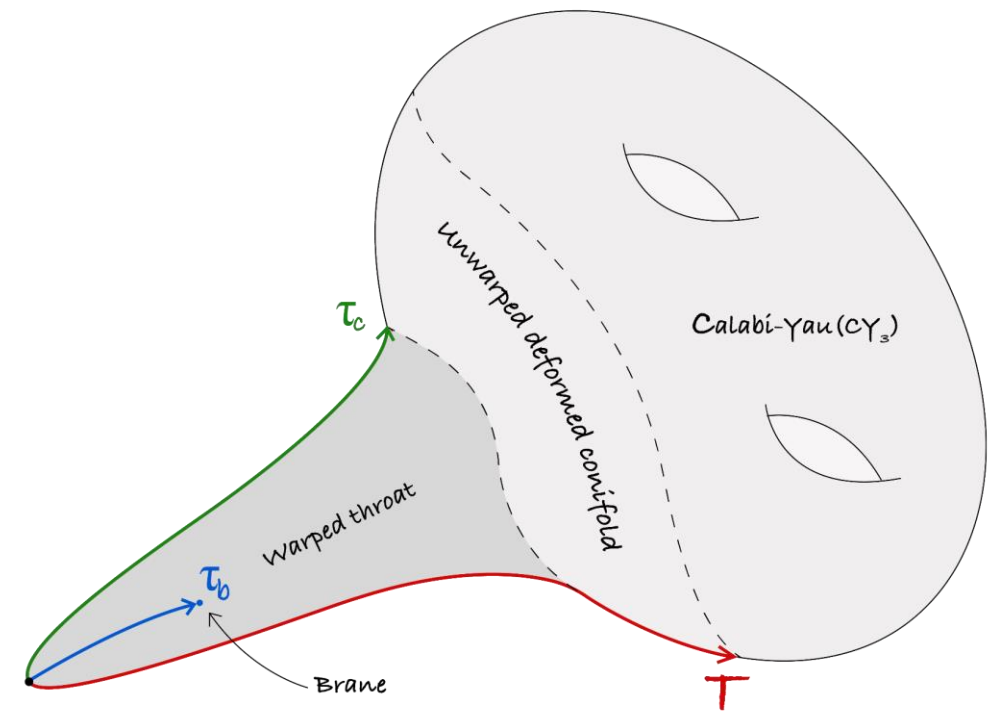
KS + Bulk

$$ds^2 = H(y)^{-1/2} g_{\mu\nu} dx^\mu dx^\nu + H(y)^{1/2} c^{1/2} g_{mn} dy^m dy^n$$



Our paper [2204.02086]

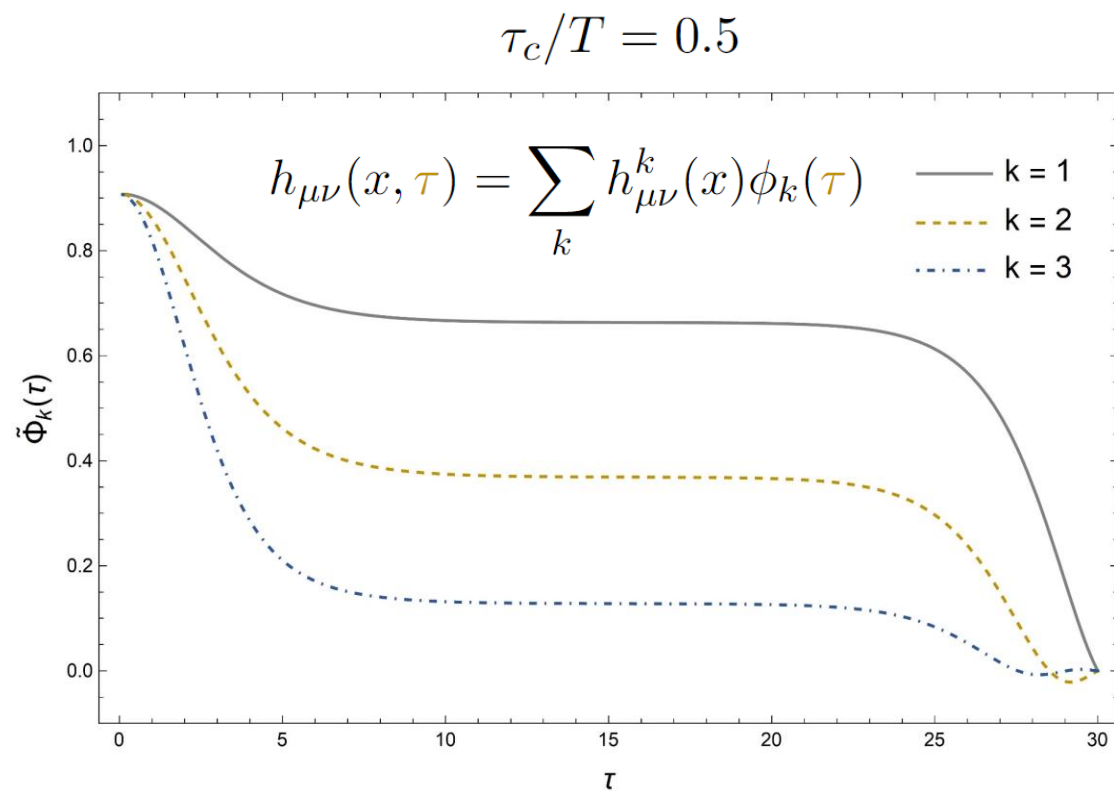
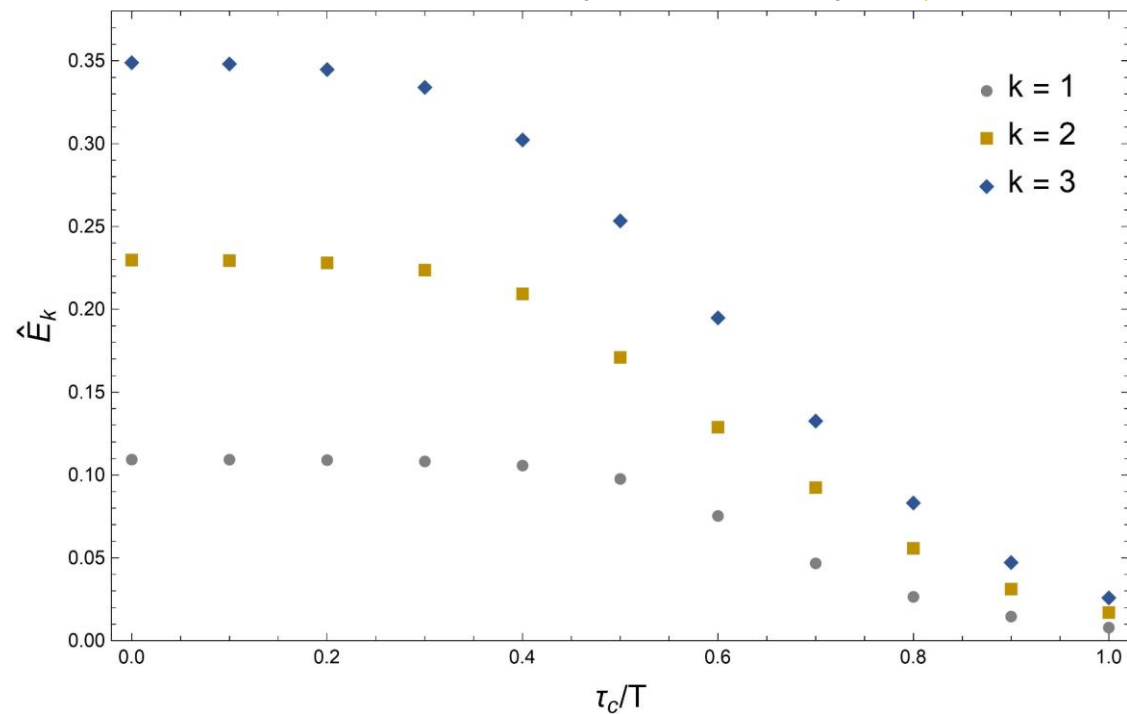
1. 4d Minkowski
2. Trivial angular solutions ($\phi_k = \phi_k(\tau)$)
3. Unwarped conifold region ($\tau_c < \tau < T$)
4. Vanishing boundary conditions on $CY_3(\tau = T)$
5. (3+1)-brane somewhere in the conifold ($\tau_b < T$)



Throat vs Bulk

There is a competition between warping and bulk size.

cf. H. Firouzjahi and S. H. H. Tye [hep-th/0512076]



Gravitational Waves

Each mode in the tower has its own wave equation

$$\square_4 h_{\mu\nu}^k - m_k^2 h_{\mu\nu}^k = T_{\mu\nu}$$

Tower of frequencies $\omega_k \sim m_k$ ($\Delta\omega = M_{KK}^w$)

	f_{GW} (Hz)	M_{KK}^w (eV)	τ_b	τ_c	$z^{1/3}$	r_{UV}	\mathcal{V}_{th}	MK
LISA	10^0	10^{-27}	195	239	1.51×10^{-47}	1.70	290	3259
LIGO-Virgo/ET	10^4	10^{-23}	168	211	1.51×10^{-43}	1.64	240	2906
UHF	10^9	10^{-18}	133	176	1.51×10^{-38}	1.57	183	2464

cf. D. Andriot and G. Lucena Gómez [1704.07392]

cf. A. K. Mishra, A. Ghosh, S. Chakraborty [2106.05558]